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结合面卸载分形模型*

缪小梅 黄筱调

(南京工业大学机械与动力工程学院,南京 210009)

摘要:基于分形接触理论以及微凸体的卸载模型,建立了结合面的分形卸载模型。在模型中考虑了弹性以及弹塑 性微凸体的卸载。通过对所建模型的数值仿真,揭示了卸载过程中接触载荷、实际接触面积、接触压力与干涉量间 的非线性关系,以及卸载过程与加载过程的不同。结果表明,结合面的卸载过程是弹性的,且卸载过程取决于加载 过程的最终状态。卸载过程中,接触面积和载荷是干涉量的函数。卸载开始时,接触面积和接触载荷都急剧减小, 并且小于加载时的值。随着卸载量的增加,接触面积很快超过了加载时的接触面积,而接触载荷在很大范围内都 远小于加载时的接触载荷。在整个干涉量区间内,接触压力都远小于卸载时的接触压力。

关键词:结合面 分形模型 卸载模型

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引言

1991年, Majumdar等研究发现大部分机械结合 面具有分形特征, 并据此提出了以分形几何为基础 的接触模型——MB模型^[1]。许多研究者在此基础 上对结合面的接触特性进行了一系列研究^[2-12]。 MB模型中仅仅考虑了微凸体的两种变形方式, 纯 弹性或纯塑性, 而对两者的过渡区间——弹塑性变 形未予考虑。单个微凸体的研究表明, 在很大的变 形量范围内, 微凸体发生的都是弹塑性变形^[13-15]。 文献[16]在此基础上提出了考虑了微凸体弹塑性 变形的结合面分形接触模型。文献[17]将结合面 的接触模型拓展到包含几何误差的情况。上述模型 只考虑了结合面在一次加载的情况, 然而对于一般 的机械零件, 通常需要承受成千上万次的载荷, 因此 卸载过程对实际的机械结合面而言是很重要的。

本文基于单个微凸体的卸载模型,提出一种结 合面的分形卸载模型。通过模型仿真研究了卸载过 程中结合面上的接触载荷,接触面积以及接触压力 的变化趋势,并将卸载过程的加载过程与文献[16] 的加载过程进行比较。

1 微凸体的卸载模型

结合面之间的接触变形实质上发生在微凸体之间,可考虑为一个半球形微凸体与刚性平面的接触问题。由于弹性变形是可逆的,发生弹性变形的微

凸体卸载与加载过程一致,满足赫兹理论。卸载过 程中微凸体接触面积 *a*_u 以及接触载荷 *f*_u 分别为

$$a_u = \pi R \delta_u \tag{1}$$

$$f_u = \frac{4E}{3} R^{1/2} \delta_u^{3/2}$$
 (2)

其中
$$\delta_{u} = \begin{cases} \delta - \delta_{m} + d & (\delta - \delta_{m} + d \ge 0) \\ 0 & (\delta - \delta_{m} + d < 0) \end{cases}$$
 (3)

微凸体 0≤ δ_u ≤δ

对于发生弹塑性变形的微凸体而言,微凸体卸 载过程与之前的受载历史相关。Johnson 首先对半 球与刚性平面接触的卸载过程进行了研究,指出即 便在加载过程中发生了大的塑性变形,其卸载过程 仍旧是完全弹性的,满足赫兹定律^[18]。Vu-Quoc^[19] 等提出了一个简单的微凸体卸载解析模型,并将模 型与有限元结果进行了比较,具有很好的一致性。 然而这个模型只适用于塑性变形比例较小的情况 下。Etsion^[20]等基于 Johnson 的假设,推导了微凸体 卸载的理论模型,同时对卸载过程进行了有限元计 算,并将两者进行了比较。Etsion 等的研究覆盖了

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作者简介:缪小梅,博士生,主要从事机床结合部特性及其精度衰退机理研究,E-mail: miaoxiaomeinjutt@126.com

弹塑性到完全塑性的全部区域。根据 Etsion 有限元 分析的结果,微凸体卸载后的残余干涉量为

$$\delta_{r} = \delta \left[1 - \frac{1}{\left(\delta/\delta_{c} \right)^{0.28}} \right] \left[1 - \frac{1}{\left(\delta/\delta_{c} \right)^{0.69}} \right] \quad (4)$$

其中
$$\delta_c = \left(\frac{\pi KH}{2E}\right)^2 R$$
 (5)

$$H = 2.8\sigma_s$$
 $K = 0.454 + 0.41\nu^{[14]}$

式中 δ。——微凸体的临界干涉量

H----材料的硬度

σ。——材料的屈服强度

K——与材料泊松比v相关的系数

卸载过程中接触载荷和接触面积与干涉量的关 系为

$$f_{u} = f\left(\frac{\delta_{u} - \delta_{r}}{\delta - \delta_{r}}\right)^{1.5\delta^{-0.0331}}$$
(6)

$$a_{u} = a \left(\frac{\delta_{u} - \delta_{r}}{\delta - \delta_{r}}\right)^{\delta^{-0.12}}$$
(7)

对于存在塑性变形的微凸体, $\delta_r < \delta_u < \delta$,即微凸体不能完全恢复成未变形前的状态,而是存在一定的残余变形

$$\delta_{u} = \begin{cases} \delta - \delta_{m} + d & (\delta - \delta_{m} + d \ge \delta_{r}) \\ \delta_{r} & (\delta - \delta_{m} + d < \delta_{r}) \end{cases}$$
(8)

加载过程的最大干涉量是影响卸载过程的主要 因素。残余干涉量与加载时的干涉量呈正比。随着 δ 的增加,残余干涉量趋近一个上限值 δ ,这时,微凸 体接近完全塑性变形,弹性恢复量相对于干涉量来 说微不足道。而当 δ 趋近于 δ_c 时,加载近似为纯弹 性的,残余干涉量趋近于零。卸载过程是完全弹性 的,但是接触载荷以及接触面积与干涉量的关系已 经不满足赫兹定律。

2 结合面的卸载模型

任意两个粗糙表面的接触都可以简化为一个粗 糙表面与一个刚性平面的接触问题,本文只考虑简 化后的情况。粗糙表面与刚性平面接触时,微凸体 接触面积 a 的分布规律与海洋面岛屿面积的分布规 律相似,分布密度服从幂函数

$$n(a) = \frac{D}{2} \frac{a_m^{\frac{D}{2}}}{a^{\frac{D}{2}+1}}$$
(9)

式中 a_m——最大接触点的面积

D——分形维数,1 < D < 2,表示轮廓高度中 高频和低频成分的分布比例,D 越大, 则高频成分所占的比例越大

根据分形理论,微凸体的半径为

$$R = \frac{a'^{D/2}}{\pi^2 G^{D-1}} \tag{10}$$

G——分形粗糙度系数,它决定了轮廓高度, G 越大,则其表面越粗糙

在弹性变形时,微凸体的截断面积为微凸体接 触面积的2倍。

根据结合面式(1)、(9)和(10),弹性微凸体卸载时总的接触面积为

$$A_{eu} = \int_{a_{ee}}^{a_{m}} n(a) a_{un} da =$$
$$\int_{a_{ee}}^{a_{m}} \frac{\pi D}{2} \frac{a_{m}^{\frac{D}{2}}}{a^{\frac{D}{2}+1}} \frac{(0.5a)^{D/2}}{\pi^{2} G^{D-1}} \delta_{un} da$$

式中 a_{ec}——弹性临界接触面积

根据文献[16],微凸体临界接触面积

$$a_{ec} = \frac{G^2}{2\left(\frac{KH}{2E}\right)^{2/(D-1)}}$$
(11)

$$a_{pc} = \frac{G^2}{110^{1/(D-1)} \left(\frac{KH}{2E}\right)^{2/(D-1)}}$$
(12)

式中 a_n——塑性临界接触面积

当 a≥a_{ee}时,微凸体发生弹性变形;当 a_{pe}≤a≤ a_{ee}时,微凸体发生弹塑性变形;当 a≤a_{pe}时微凸体发 生塑性变形。

根据分形理论,微凸体的高度

$$\delta = G^{D-1} a'^{(2-D)/2} \tag{13}$$

对于弹塑性变形时接触面积与截断面积的关 系,本文采用与 KE 模型相同形式的函数来拟 合^[14],同时保持函数的连续性,即在弹性临界变形 处接触面积为截断面积的一半,塑性临界变形接触 面积等于截断面积

$$a' = a \left(\frac{a}{a_{pc}}\right)^{\frac{1}{D-1} \lg^{10} - \lg^{2}}$$
(14)

根据式(4)、(7)和(9),弹塑性微凸体卸载时总的接触面积

$$A_{epu} = \int_{a_{pc}}^{a_{ec}} n(a) a da =$$

$$\int_{a_{pc}}^{a_{ec}} \frac{D}{2} \frac{a_{m}^{\frac{D}{2}}}{a^{\frac{D}{2}}} \left(\frac{\delta_{u} - \delta_{r}}{\delta - \delta_{r}}\right)^{\delta^{-0.12}} da \qquad (15)$$

总的接触面积

$$A_u = A_{eu} + A_{epu} \tag{16}$$

由式(3)、(4)可知, δ_{u} 和 δ_{r} 是加载时结合面最 大接触点的面积 a_{m} 、加载时微凸体的最大接触面积 a和干涉量d的函数。根据式(13)、(14), δ 是加载 时结合面最大接触点的面积 a_{m} 、加载时微凸体的最 大接触面积a和干涉量d的函数。因此总的接触面 积是干涉量d以及加载时最大接触点面积 a_{m} 的函 数,也就是说结合面卸载时的接触面积与干涉量的 关系只取决于结合面加载时的最终状态。根据 式(2)、(9)、(10)和(13),卸载时发生弹性变形的 微凸体上总的接触载荷为

$$F_{eu} = \int_{a_{ec}}^{a_{m}} n(a) f_{u} da = \int_{a_{ec}}^{a_{m}} \frac{2DE}{3} \frac{a_{m}^{\frac{D}{2}} (0.5a)^{D/4}}{\pi G^{0.5(D-1)} a^{\frac{D}{2}}} \delta_{u}^{1.5} da$$
(17)

根据文献[16],加载时发生弹塑性变形的微凸体的载荷与接触面积的关系为

其中

$$f = 1.6156\sigma_s a \left(\frac{a}{a_{ec}}\right)^c$$
 (18)
其中
 $C = \frac{\lg 1.733}{\lg 2 - \lg \left(110^{\frac{1}{D-1}}\right)}$

根据式(4)、(6)、(9)和(18),卸载时发生弹塑性变形的微凸体上总的接触载荷为

$$F_{epu} = \int_{a_{pc}}^{a_{ec}} n(a) f_u da =$$
$$\int_{a_{pc}}^{a_{ec}} \frac{D}{2} \frac{1.615}{a} \frac{6\sigma_s a_m^{\frac{D}{2}} a}{a^{\frac{D}{2}+1}} \left(\frac{a}{a_{ec}}\right)^C \left(\frac{\delta_u - \delta_r}{\delta - \delta_r}\right)^{1.5\delta^{-0.0331}} da$$
(19)

同样地,总的接触载荷是干涉量 d 以及加载时最大接触点面积 a_m 的函数,结合面卸载时的接触面积与 干涉量的关系只取决于结合面加载时的最终状态。

3 数值仿真

仿真参数:结合面材料为 Q235, E = 230 GPa, $\nu = 0.3$, $\sigma_s = 235$ MPa, $G = 10^{-14}$ m, D = 1.8, 加载最 终实际接触面积 $A_R = 0.01A_0$, A_0 为名义接触面积, 这里取 $A_0 = 1000$ mm²。







图 1 为结合面加载和卸载接触面积,接触面积 是干涉量的非线性函数。由图可见,在刚开始卸载 时,接触面积小于加载时的接触面积。这是由于在 刚开始卸载时,原本发生塑性以及弹塑性接触的微 凸体,由于存在残余变形,很快脱离接触。但是随着 卸载量的增加,卸载时的接触面积很快就超过了加 载时的接触面积。

图 2 为结合面加载和卸载时的接触载荷,接触

载荷是干涉量的非线性函数。由图2可见,卸载一 开始,接触载荷急剧减少,变化的速率远大于加载时 的变化率。产生这一现象的原因是在卸载过程中, 原本发生弹塑性以及塑性变形的微凸体上的载荷的 迅速减少。但是在接近卸载完成时,卸载时的接触 载荷略大于加载时的接触载荷。



图 3 为结合面加载和卸载时实际接触面积上的 平均接触压力。显然在整个卸载过程中,卸载的接 触压力远小于加载时的接触压力。这是由于加载时 结合面既发生弹性变形也发生塑性变形,而结合面 的卸载过程是完全弹性的。显然,塑性变形相对弹 性变形具有更好的承载能力。



4 结论

(1) 微凸体的卸载过程完全弹性,即便加载过 程是包含塑性的,但是已经不服从赫兹接触理论。 微凸体的卸载过程取决于微凸体加载时的最终状态。结合面卸载后,存在塑性变形的微凸体不能完 全恢复,从而使得结合面存在残余变形。

(2)结合面的卸载过程是完全弹性的,再次加载时只要干涉量不超过上次加载时的干涉量,再次加载过程将重复卸载的过程。结合面的卸载过程取决于结合面加载时的最终状态。

(3)卸载时,结合面的接触载荷、接触面积和接 触压力是加载时结合面最大接触点面积和卸载时结 合面间干涉量的函数。结合面的卸载过程取决于结 合面加载时的最终状态。

(4) 在刚开始卸载时,结合面接触面积和接触

载荷迅速减少,远小于加载时的接触面积和接触载 荷。随着卸载过程的继续,卸载的接触面积很快超 过加载时的接触面积。但是接触载荷直到快要卸载

完成时才略大于加载时的接触载荷。在整个卸载过 程中结合面上的接触压力远小于加载时的接触压 力。

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Online Identification and Correction of Orifice Inlet Edge Bluntness Based on Downstream Flow Field Distortion

Li Hongwen^{1,2} Zhang Tao¹

(1. School of Electrical Engineering and Automation, Tianjin University, Tianjin 300072, China

2. School Office, Northeast Petroleum University, Daqing 163318, China)

Abstract: Online identification of orifice bluntness was put forward, and specific solutions was given. Orifice discharge coefficient of CFD simulation on normal condition and bluntness conditions were compared with GB/T 21446—2008 to verify the validity of the simulation. To achieve the identification, in the rear of the flange pressure points (P_1, P_2) another pressure point P_3 was installed and the differential pressure ratio factor (η) was calculated. The factor reflected the degree of distortion of orifice plate flow field relative to the normal condition, which was the ratio factor indirectly reflected the degree of bluntness. The best position of pressure points P_3 was determined by the analysis of simulation results, and the experience formula about the ratio factor with bluntness was further concluded. Through the formula, online identification of bluntness could be realized. The effectiveness of the identification algorithm was verified through real flow experiments. The error of calculated value was within $\pm 1.0\%$ for sharpness correction coefficient (b_k) , which proved that the algorithm was practicable in engineering. **Key words**: Orifice meter Inlet edge bluntness Online identification Differential pressure ratio factor

Flow field distortion

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Fractal Unloading Model of Joint Interfaces

Miao Xiaomei Huang Xiaodiao

(School of Mechanical and Power Engineering, Nanjing University of Technology, Nanjing 210009, China)

Abstract: A fractal unloading model of joint interfaces was established based on fractal contact theory and unloading model of asperity. The unloading of asperities both elastic and elastic-plastic deformation was taken into account in the model. By numerical simulation, the nonlinear relationship among the real contact areas, contact force, contact pressure and interference was obtained, as was as the difference between loading and unloading. The result showed that the unloading process of joint interfaces was elastic, and depended on the final condition of loading. During the unloading process, the real contact area and contact force depended on the interference of joint interface. When the loading process started, both the real contact area and contact force sharply decreased and smaller than the values during loading. As the interference increased, the real contact area became larger than the value during loading soon, while the contact force was smaller than the value during loading in a wide range of interference. The unloading contact pressure was much less than the loading contact pressure within the entire interference interval.

Key words: Joint interfaces Fractal modeling Unloading modeling